

Quantum Entanglement Lecture 3 2006-10-09

*Quantum Mechanics is calculation of probabilities
sigma matrices
observables, eigenvalues and eigenvectors are orthogonal
can measure component of electron spin in any x,y,z
example, unit pointer in any arbitrary direction*

Prof. Leonard Susskind; videos on [Stanford on iTunes U](#)
[Susskind's Blog: Physics for Everyone](#)

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01:00 *Quantum Mechanics is calculation of probabilities*

$$|a\rangle = \alpha |up\rangle + \beta |down\rangle$$

$$\alpha \beta \text{ are probabilities } \alpha^2 + \beta^2 = 1$$

Hermitean matrices, some identities

- Transpose and complex conjugate, $M_{ij} = M_{ji}^*$
- M^\dagger (dagger \dagger is Hermitian conjugate)
- Real elements on diagonal,
- Complex conjugate off-diagonal

$$\langle b|a\rangle = \langle a|b\rangle^* = [b1 \quad b2] \begin{bmatrix} a1 \\ a2 \end{bmatrix} = \left([a1 \quad a2] \begin{bmatrix} b1 \\ b2 \end{bmatrix} \right)^*$$

$$\langle b|a\rangle = b1^*a1 + b2^*a2$$

$$\langle a|b\rangle = a1^*b1 + a2^*b2$$

$$\langle a|b\rangle^* = b1^*a1 + b2^*a2 = \langle b|a\rangle$$

$\langle bra|$ is row vector, complex conjugate; $|ket\rangle$ is normal column vector

Hermitian calculation:

$$\langle b|M|a\rangle = \langle a|M^\dagger|b\rangle^* = [b1 \quad b2] \begin{bmatrix} m11 & m12 \\ m12^* & m22 \end{bmatrix} \begin{bmatrix} a1 \\ a2 \end{bmatrix} = \left([a1 \quad a2] \begin{bmatrix} m11 & m12 \\ m12^* & m22 \end{bmatrix} \begin{bmatrix} b1 \\ b2 \end{bmatrix} \right)^*$$

could multiply this out but as we know $M = M^\dagger$ and above result it appears obvious.

- a is real
- expectation value of a

36:00 any observable can be represented as a collection of real numbers

39:00 *eigenvalues and eigenvectors*

$$M|a\rangle = \lambda_a |a\rangle$$

- If M is a Hermitean matrix then λ_a is real
- Take inner product
 - o $\langle a|M|a\rangle = \lambda_a \langle a|a\rangle$ (λ_a is a number)
 - o $\langle a|M|a\rangle$ for M =hermitean $\rightarrow a$ is real
 - o Therefore λ_a is real

- If M is an observable then the value λ_a is a measurable value.

44:00

The eigenvector (collection of λ eigenvalues) is the state of the system.

54:00 *sigma matrices*

sigma 3 z-axis (sometimes called the spin operator)

$$\sigma_3 = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

eigenvector $\{1 \ 0\}, \{0 \ -1\}$, eigenvalues $+1, -1$

$$\sigma_3 |1 \ 0\rangle = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = +1 \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$\sigma_3 |0 \ -1\rangle = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{pmatrix} 0 \\ -1 \end{pmatrix} = -1 \begin{pmatrix} 0 \\ -1 \end{pmatrix}$$

sigma 1 x-axis

$$\sigma_1 = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

eigenvector $\{1 \ 1\}, \{1 \ -1\}$, eigenvalues $+1, -1$ (divided by $\sqrt{2}$)

$$\sigma_1 |1 \ 1\rangle = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

$$\sigma_1 |1 \ -1\rangle = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = \begin{pmatrix} -1 \\ 1 \end{pmatrix} = \frac{-1}{\sqrt{2}} \begin{pmatrix} -1 \\ 1 \end{pmatrix}$$

sigma 2 y-axis

$$\sigma_2 = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$$

eigenvector $\{1 \ i\}, \{1 \ -i\}$, eigenvalues $+1, -1$

$$\sigma_2 |1 \ i\rangle = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{pmatrix} 1 \\ i \end{pmatrix} = +1 \begin{pmatrix} 1 \\ i \end{pmatrix}$$

$$\sigma_2 |1 \ -i\rangle = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{pmatrix} 1 \\ -i \end{pmatrix} = -1 \begin{pmatrix} 1 \\ -i \end{pmatrix}$$

sigma identities

$$\sigma_1^2 = \sigma_3^2 = \sigma_2^2 = I$$

$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$$\sigma_1 \sigma_2 = -i \sigma_3 \quad \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} = -i \begin{bmatrix} i & 0 \\ 0 & -i \end{bmatrix}$$

$$\sigma_3 \sigma_1 = -i \sigma_2 \quad \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} = -i \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}$$

$$\sigma_2 \sigma_3 = -i \sigma_1 \quad \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} = -i \begin{bmatrix} 0 & i \\ i & 0 \end{bmatrix}$$

note if you reverse the sigma matrix multiply order, the sign is reversed:

$$\sigma_2 \sigma_1 = +i \sigma_3$$

$$\sigma_1 \sigma_3 = +i \sigma_2$$

$$\sigma_3 \sigma_2 = +i \sigma_1$$

60:00 *observables, eigenvalues and eigenvectors are orthogonal*

suppose an observable with 2 eigenvalue
(normally 3x3 M has 3, 4x4 has 4, ...)

an observable with 2 eigenvalue

$$M|a\rangle = \lambda_a |a\rangle \quad \text{eigenvector } a \text{ with eigenvalue } \lambda_a$$

$$M|b\rangle = \lambda_b |b\rangle \quad \text{eigenvector } b \text{ with eigenvalue } \lambda_b$$

if λ_a & λ_b are different then they are orthogonal
(the vectors a,b are the states of the systems)

proof

$$1 \quad \langle b|M|a\rangle = \lambda_a \langle b|a\rangle \quad \text{multiply by } b$$

$$2 \quad \langle a|M|b\rangle = \lambda_b \langle a|b\rangle \quad \text{multiply by } a$$

take complex conjugate of 2

$$2a \quad \langle a|M|b\rangle^* = \lambda_b^* \langle a|b\rangle^* \quad \text{or}$$

$$2b \quad \langle b|M|a\rangle = \lambda_b \langle b|a\rangle$$

subtract 2b from 1 giving:

$$3 \quad 0 = (\lambda_a - \lambda_b) \langle b|a\rangle$$

if $\lambda_a \neq \lambda_b$, then a is orthogonal to b

--- exercise: check that σ vectors are orthogonal

use σ_2 defn. and identities:

$$\sigma_2 = \sigma_2^\dagger; \lambda_a = +1, \lambda_b = -1$$

$$|a\rangle = \begin{pmatrix} 1 \\ i \end{pmatrix} \quad \langle a| = (1 \quad -i) \quad |b\rangle = \begin{pmatrix} 1 \\ -i \end{pmatrix} \quad \langle b| = (1 \quad i)$$

$$4 \quad \langle b|\sigma_2|a\rangle = \lambda_a \langle b|a\rangle \quad (1 \quad i) \begin{bmatrix} 0 & i \\ -i & 0 \end{bmatrix} \begin{pmatrix} 1 \\ i \end{pmatrix} = +1(1 \quad i) \begin{pmatrix} 1 \\ i \end{pmatrix}$$

$$5 \quad \langle a|\sigma_2|b\rangle = \lambda_b \langle a|b\rangle \quad (1 \quad -i) \begin{bmatrix} 0 & i \\ -i & 0 \end{bmatrix} \begin{pmatrix} 1 \\ -i \end{pmatrix} = -1(1 \quad -i) \begin{pmatrix} 1 \\ -i \end{pmatrix}$$

take complex conjugate of 5

$$5a \quad \langle a|\sigma_2|b\rangle^* = \lambda_b^* \langle a|b\rangle^*$$

$$5b \quad \langle a|\sigma_2|b\rangle^* = (1 \quad +i) \begin{bmatrix} 0 & i \\ -i & 0 \end{bmatrix} \begin{pmatrix} 1 \\ +i \end{pmatrix} = \langle b|\sigma_2|a\rangle$$

$$5c \quad \lambda_b^* \langle a|b\rangle^* = -1(1 \quad +i) \begin{pmatrix} 1 \\ i \end{pmatrix} = \lambda_b \langle b|a\rangle$$

$$5d \quad \langle b|\sigma_2|a\rangle = \lambda_b \langle b|a\rangle$$

subtracting 4-5d; as $\lambda_a \neq \lambda_b$ then $\langle b|a\rangle = 0$

$$6 \quad 0 = \lambda_a \langle b|a\rangle - \lambda_b \langle b|a\rangle$$

notes:

1 as any $\sigma = \sigma^\dagger$, and a,b can be replaced with any σ eigenvector the above is sufficient for σ_1, σ_3

2 could have used (3) and just calculate $\langle b|a\rangle = \langle a|b\rangle^* = 0$

72:00 can measure component of electron spin in any x,y,z

- but not simultaneously
- sigma vectors are the x, y, z components

let:

M be a system of observables; with λ_a , the probability of M being in state a (eigenvector a)

if you prepare the system as b – what is the probability that will be in the state a?

the probability is the square of the dot product of b and a

$$\langle a|b\rangle \langle a|b\rangle^*$$

dot product of a,b times the dot product of the conjugate (square)

notes:

1: a,b are unit vectors;

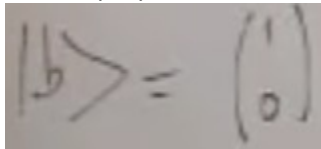
2: $\langle a|b\rangle$ is a complex number, the square (complex conjugate) is a real number;

3: the probability, a real number, will always be ≤ 1 ;

4: if a,b are orthogonal the probability $\langle a|b\rangle$ is zero

80:00 if a,b were real vectors then the dot product squared $\langle a|b\rangle \langle a|b\rangle$ would be the cosine between them. But a,b are in general, complex so that the square must be the complex conjugate.

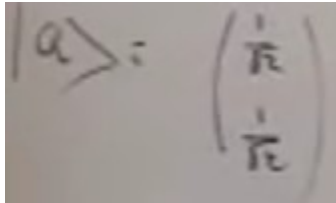
82:00 prepare as σ_3 , test as σ_1 – probability is $\frac{1}{2}$


$$|b\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

prepared as σ_3 , which means $\sigma_3 = +1$

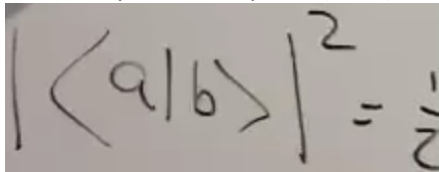
this is done by placing a magnetic field in the “up” or “+z” direction. The observable, the spin, will be +1

now we will measure σ_1 (horizontal “+x” direction) and ask what is the probability we will get +1?


$$|a\rangle = \begin{pmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{pmatrix}$$

a is an eigenvector of σ_1

the inner product $\langle a|b\rangle$ is $(1/\sqrt{2})$, the amplitude is $(1/\sqrt{2})(1/\sqrt{2}) = (1/2)$


$$|\langle a|b\rangle|^2 = \frac{1}{2}$$

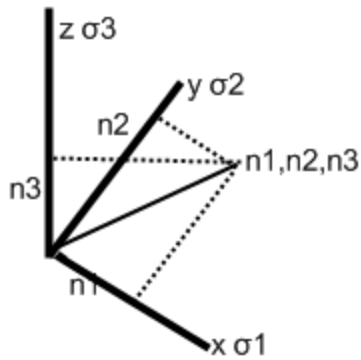
the form $|\langle a|b\rangle|^2$ is convention for $\langle a|b\rangle \langle a|b\rangle^*$

similar probability for “-x”; $(-1/\sqrt{2})(-1/\sqrt{2}) = (1/2)$

86:00 if line up spin in one direction, measure in another – probability is $\frac{1}{2}$

93:00 example, unit pointer in any arbitrary direction.

--- always get +1 or -1 probabilities ...



vector n (n1,n2,n3) projected onto sigma 3-D axis

n1 n2 n3 pointer components (unit normalized)
 components of the spin are:

$\sigma \cdot n$ dot product

$$\sigma_1 * n_1 + \sigma_2 * n_2 + \sigma_3 * n_3$$

multiplying each sigma matrix by n, then adding:

$$\begin{bmatrix} 0 & n_1 \\ n_1 & 0 \end{bmatrix} + \begin{bmatrix} 0 & -in_2 \\ in_2 & 0 \end{bmatrix} + \begin{bmatrix} n_3 & 0 \\ 0 & -n_3 \end{bmatrix} = \begin{bmatrix} n_3 & n_1 - in_2 \\ n_1 + in_2 & -n_3 \end{bmatrix}$$

we get a Hermitian matrix:

$$\sigma \cdot n = \sigma_1 * n_1 + \sigma_2 * n_2 + \sigma_3 * n_3 = \begin{bmatrix} n_3 & n_1 - in_2 \\ n_1 + in_2 & -n_3 \end{bmatrix}$$

let us square the dot product; $(\sigma \cdot n) * (\sigma \cdot n)$ to show this equals one:

$$(\sigma_1 * n_1 + \sigma_2 * n_2 + \sigma_3 * n_3) * (\sigma_1 * n_1 + \sigma_2 * n_2 + \sigma_3 * n_3)$$

the normal terms equal one:

$$\sigma_1 * n_1 * \sigma_1 * n_1 + \sigma_2 * n_2 * \sigma_2 * n_2 + \sigma_3 * n_3 * \sigma_3 * n_3$$

grouping,

$(n_1 * n_1 * \sigma_1 * \sigma_1) +$ we know that any sigma squared = one, leaves us with:

$$n_1 * n_1 + n_2 * n_2 + n_3 * n_3 = 1 \quad \text{as vector n is unit normalized by defn.}$$

the cross product terms all disappear:

$$\sigma_1 * n_1 * \sigma_2 * n_2 + \sigma_1 * n_1 * \sigma_3 * n_3 +$$

$$\sigma_2 * n_2 * \sigma_1 * n_1 + \sigma_2 * n_2 * \sigma_3 * n_3 +$$

$$\sigma_3 * n_3 * \sigma_1 * n_1 + \sigma_3 * n_3 * \sigma_2 * n_2$$

re-grouping:

$$\sigma_1 * n_1 * \sigma_2 * n_2 + \sigma_2 * n_2 * \sigma_1 * n_1 +$$

$$\sigma_1 * n_1 * \sigma_3 * n_3 + \sigma_3 * n_3 * \sigma_1 * n_1 +$$

$$\sigma_2 * n_2 * \sigma_3 * n_3 + \sigma_3 * n_3 * \sigma_2 * n_2$$

note that each line is equal to zero because of reverse order of matrix multiplication:

$$\sigma_1 * n_1 * \sigma_2 * n_2 = \begin{bmatrix} 0 & n_1 \\ n_1 & 0 \end{bmatrix} \begin{bmatrix} 0 & -in_2 \\ in_2 & 0 \end{bmatrix} = \begin{bmatrix} n_1 + in_2 & 0 \\ 0 & n_1 - in_2 \end{bmatrix}$$

$$\sigma_2 * n_2 * \sigma_1 * n_1 = \begin{bmatrix} 0 & -in_2 \\ in_2 & 0 \end{bmatrix} \begin{bmatrix} 0 & n_1 \\ n_1 & 0 \end{bmatrix} = \begin{bmatrix} n_1 - in_2 & 0 \\ 0 & n_1 + in_2 \end{bmatrix}$$

or $(\sigma_1 * n_1 * \sigma_2 * n_2) = -(\sigma_2 * n_2 * \sigma_1 * n_1)$

the cross products equalling zero can be deduced by the sigma identities above.